Heat Conduction in Ceramic Coatings: Relationship between Microstructure and Effective Thermal Conductivity

Technical report for Task 1 (period of performance Nov. 7, 1996 - Feb. 6, 1997)

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Research on the effective thermal conductivity of ceramic coatings and its relation to the microstructure continued. Experimental data on thermal conductivity of ceramic coatings provided by NASA (paper by R. Miller et al "The effect of plasma spray parameters on density, cyclic life and thermal conductivity of a ZrO₂-Y₂O₃ thermal barrier coatings", Presented at the ASM-AIME/ Materials Week, Cleveland, Ohio, October 1995) indicates the necessity of construction a theoretical framework for the analysis of anisotropic conductivity in terms of the microstructural parameters. Our effort was focused on this task.

We employed the approach similar to the one developed earlier for effective elastic properties of materials with defects and developed it further for the thermal conductivity problem. We started with the analysis of the influence of one pore on the overall heat flux. This influence is strongly dependent on the shape of the pore. We model pores as full insulators (this modelling can be viewed as the first approximation and can later be refined to account for a finite, albeit small, conductivity of a pore).

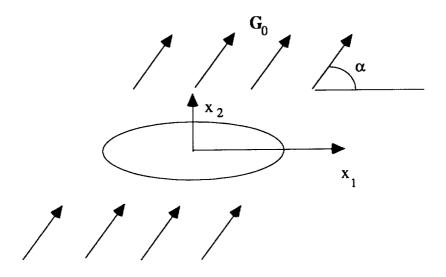
We first analyzed pores of *elliptical* shapes. Such shapes, although simple, constitute an important example: they demonstrate the importance of shape factors. Also, they cover both cracks and circles as limiting cases, and allow one to consider the practically important case of mixtures of cracks/pores. We start with the analysis of an influence of one such pore on the overall heat flux. A summary of our findings is as follows.

One elliptical hole. We consider a 2-D infinite medium of thermal conductivity k_0 with an insulating (k=0) elliptical inclusion (axes 2a and 2b). The heat flux vector $\mathbf{Q}_0 = k_0 \mathbf{G}_0$ far away from the inclusion is inclined at an angle α to the major semi-axis of the ellipse.

The scalar potential Ω on the contour of the elliptical inclusion is

$$\Omega = -k_0 G_0(a+b)\cos(\theta+\alpha) \tag{1}$$

where θ is angular elliptical coordinate of the point on the boundary of an inclusion, as given by conformal mapping $z = R(m\xi + 1/\xi)$, $\xi = e^{i\theta}$, R = (a+b)/2, m = (a-b)/(a+b).



The macroscopic heat flux for the representative area element (RAE) containing a hole is:

$$Q = \frac{1}{A} \int_{\Gamma} \Omega \mathbf{n} \, d\Gamma = \frac{1}{A} \int_{A} Q \, dA - \frac{1}{A} \int_{\gamma} \Omega \mathbf{n} \, d\Gamma = Q_0 + \Delta Q \tag{2}$$

where A- representative area, γ - hole boundary, A_s - area of the solid phase and $\Omega(G_1, G_2)$ is a scalar potential such that $Q_1 = \frac{\partial \Omega}{\partial G_1}$, $Q_2 = \frac{\partial \Omega}{\partial G_2}$.

In the case of the elliptical hole, our calculations yield

$$n_{1} = \frac{(m-1)\cos\theta}{\sqrt{1+m^{2}-2m\cos2\theta}}, \quad n_{2} = \frac{(m+1)\sin\theta}{\sqrt{1+m^{2}-2m\cos2\theta}}, \quad d\Gamma = R\sqrt{1+m^{2}-2m\cos2\theta} \ d\theta$$
(3)

$$\Delta Q_1 = -\frac{k_0 G_0}{A} \pi b(a+b)\cos \alpha = -\frac{k_0}{A} \pi b(a+b)G_1 \tag{4}$$

$$\Delta Q_2 = -\frac{k_0 G_0}{A} \pi b(a+b) \sin \alpha = -\frac{k_0}{A} \pi b(a+b) G_2$$
 (5)

Thus, the heat flux change ΔQ can be represented as $\Delta Q = H \cdot G$ where H is the following second rank tensor:

$$H = -\frac{k_0}{A}\pi(a+b)(be_1e_1 + ae_2e_2) = -\frac{k_0}{A}\pi(abI + a^2e_2e_2 + b^2e_1e_1)$$
 (6)

The additional (due to the hole) heat energy potential $\Delta\Omega$ is expressed in terms of tensor H:

$$\Delta\Omega = \frac{1}{2}\mathbf{G} \cdot \mathbf{H} \cdot \mathbf{G} \tag{7}$$

In the case of many holes, we first consider the approximation of non-interacting holes: we assume that each hole experiences the influence of the same far-field temperature gradient G unperturbed by the presence of the other holes. Then,

$$H = \sum H^{(i)} = -\frac{k_0}{A} \pi \left[\sum a^{(i)} b^{(i)} I + \sum (a^2 n n + b^2 m m)^{(i)} \right] = -k_0 (pI + \beta)$$
 (8)

where $m^{(i)}$ and $n^{(i)}$ are the unit vectors aligned with the semi-axes $(a^{(i)})$ and $b^{(i)}$, correspondingly) of the *i*-th hole; H is expressed in terms of scalar porosity $p = (1/A)\pi \sum a^{(i)}b^{(i)}$ and a symmetric second rank tensor $\beta = (1/A)\pi \sum (a^2nn + b^2mm)^{(i)}$.

This establishes the basic framework for the analysis of various practically important distributions of shapes and orientations of pores. This work is currently underway.

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